



The modulated antiferromagnetic structures in multiferroic FeVO₄: A ⁵⁷Fe Mössbauer spectroscopy investigation

D. Colson, A. Forget, Pierre Bonville

► To cite this version:

D. Colson, A. Forget, Pierre Bonville. The modulated antiferromagnetic structures in multiferroic FeVO₄: A ⁵⁷Fe Mössbauer spectroscopy investigation. *Journal of Magnetism and Magnetic Materials*, 2015, 378, pp.529-534. 10.1016/j.jmmm.2014.11.078 . cea-01366305

HAL Id: cea-01366305

<https://hal-cea.archives-ouvertes.fr/cea-01366305>

Submitted on 14 Sep 2016

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



The modulated antiferromagnetic structures in multiferroic FeVO₄: A ⁵⁷Fe Mössbauer spectroscopy investigation

D. Colson, A. Forget, P. Bonville*

CEA, Centre de Saclay, DSM/IRAMIS/Service de Physique de l'Etat Condensé, 91191 Gif-sur-Yvette Cedex, France



ARTICLE INFO

Article history:

Received 3 September 2014

Received in revised form

21 November 2014

Accepted 26 November 2014

Available online 29 November 2014

Keywords:

Multiferroics

Mössbauer spectroscopy

Modulated magnetic structure

FeVO₄

ABSTRACT

We present an investigation of the two incommensurate magnetic phases in the multiferroic material FeVO₄ using magnetic susceptibility and ⁵⁷Fe Mössbauer spectroscopy measurements. The susceptibility shows anomalies at the phase transitions at 15.7 and 23 K. We devise Mössbauer lineshapes appropriate for planar elliptical and collinear modulated magnetic structures and show that they reproduce very well the Mössbauer spectra in FeVO₄, in full qualitative agreement with a previous neutron diffraction study. Quantitatively, our spectra provide precise determinations of the characteristics of the elliptical and modulated structures which match rather well the neutron diffraction results. We find that the hyperfine field elliptical modulation persists as $T \rightarrow 0$, which we attribute to an anisotropy of the hyperfine interaction since a moment modulation is forbidden at $T=0$ for a spin only ion like Fe³⁺.

© 2014 Elsevier B.V. All rights reserved.

1. Introduction

It is now generally accepted that most multiferroic materials, i.e. materials where magnetic and electric dipole moments are long range ordered and coupled [1], are associated with non-collinear spin density waves (SDW) incommensurate with the lattice [2,3], like cycloidal or spiral arrangements. The weak coupling case is illustrated by BiFeO₃, where ferroelectric order [4] ($T_c=1143$ K) takes place at a much higher temperature than the antiferromagnetic (AF) order ($T_N=643$ K), which consists in an incommensurate cycloidal moment arrangement [5,6]. More recently, a new class of multiferroics has been discovered, pertaining to the strong coupling case where ferroelectricity is induced by the non-collinear SDW [7,8] and appears therefore simultaneously with the SDW order. Examples of this class are TbMnO₃ [7] and TbMn₂O₅ [9]. The link between ferroelectricity and non-collinear magnetic order can be obtained in a continuum theory by considering the so-called Lifshitz invariant coupling the electric polarisation **P** and the gradient $\nabla \mathbf{M}$ of the inhomogeneous magnetisation [2]. The spontaneous polarisation can be viewed as due to an equivalent polarising electric field $\mathbf{E}_U = \gamma[(\mathbf{M} \cdot \nabla) \mathbf{M} - \mathbf{M}(\nabla \cdot \mathbf{M})]$, where γ is the coupling parameter. The volume averaged polarisation can then be shown to be non-zero for spiral (elliptic) structures and to vanish for collinear SDW.

Helical order usually appears in AF materials as a result of exchange frustration, when for instance first and second neighbour

exchange interactions are of the same magnitude [10]. Therefore, the strong coupling in AF multiferroics can be expected to be magnetically frustrated, which is reflected in the fact that the actual ordering temperature (T_N) is much lower than the exchange coupling, whose magnitude is the paramagnetic Curie temperature θ_p . For ferromagnetic interactions, helical structures are induced by antisymmetric (or Dzyaloshinski–Moriya) exchange [11]. Neutron diffraction is by far the best method allowing observation and characterisation of these incommensurate magnetic structures, but Mössbauer spectroscopy, although being a local technique which does not give access to the propagation vector, can be rather selective through lineshape analysis, especially in the case of amplitude modulated structures.

We report here on a detailed ⁵⁷Fe Mössbauer spectroscopy study of the strong coupling multiferroic FeVO₄, which presents the property of showing two magnetic transitions towards incommensurate phases [12–14] like TbMnO₃ [15]. The high temperature magnetic phase ($15.7 \text{ K} < T < 23 \text{ K}$, phase I) is a collinear sine-wave modulated structure which is not ferroelectric. The low temperature phase ($T < 15.7 \text{ K}$, phase II) is a planar non-collinear elliptical structure showing a spontaneous electric polarisation. The major elliptical axis in phase II coincides with the moment direction in phase I [13]. Early Mössbauer spectra have been reported in this compound [16,17], but they could not be thoroughly interpreted due to the lack of knowledge of the magnetic structure. We show here that the peculiar shapes of the Mössbauer spectra in both magnetic phases are entirely compatible with the magnetic structures determined by neutron diffraction [13]. We

* Corresponding author.

E-mail address: pierre.bonville@cea.fr (P. Bonville).

also present the thermal variation of the characteristics of the magnetic structures.

2. Sample synthesis and magnetic characterisation

The polycrystalline FeVO_4 sample was synthesized by heating a 1:1 molar mixture of V_2O_5 and Fe_2O_3 (hematite) at 550, 625, 700 and 715 °C during 10 h at each temperature and with intermediate grindings. FeVO_4 crystallises in the $\text{P}\bar{1}$ space group and the triclinic unit cell contains 3 different crystallographic Fe sites with very low point symmetry (inversion $\bar{1}$) [17]. All the diffraction peaks of the XRD pattern of our sample can be indexed based on the parameters determined from a single crystal in Ref. [12]. No trace of impurity phases is observed, as can be seen in Fig. 1.

The magnetic susceptibility χ of FeVO_4 was measured in a field of 0.1 T between 2 and 25 K, and with a field of 1 T from 25 K up to room temperature, using a Cryogenic Vibrating Sample Magnetometer. Down to about 100 K, the inverse susceptibility (Fig. 2) follows a Curie–Weiss law $\chi = \mu_{\text{eff}}^2 / 3k_B(T - \theta_p)$ with $\mu_{\text{eff}} = 5.93 \mu_B$ (very close to the effective moment $5.916 \mu_B$ expected for the $S=5/2$ ion Fe^{3+}) and $\theta_p \simeq -95$ K, characteristic of antiferromagnetic interactions. On further lowering the temperature, an anomaly occurs at $T_{N1} \simeq 23$ K, marking the onset of the first transition, and an inflexion point is observed at $T_{N2} = 15.7$ K, signalling the second transition. A remarkable feature is that the ratio $|\theta_p|/T_N$ is close to 5, which indicates a rather high degree of frustration of the magnetic interactions, as already observed in Refs. [13,18]. The deviation of $1/\chi$ from the Curie–Weiss law below 100 K, which cannot be attributed to crystal electric field effects absent for Fe^{3+} , is in line with this picture and show the persistence of strong short range spin correlations far in the paramagnetic phase [18].

Before describing the Mössbauer data, in the two next sections we first recall the effect on the spectra of incommensurate equal moment structures, then we compute the unusual lineshapes associated with incommensurate modulated magnetic structures, either planar elliptical or collinear.

3. Spectral effects for an incommensurate equal moment spiral or cycloidal arrangement

For the $L=0, S=5/2$ Fe^{3+} ion in the magnetically ordered phase

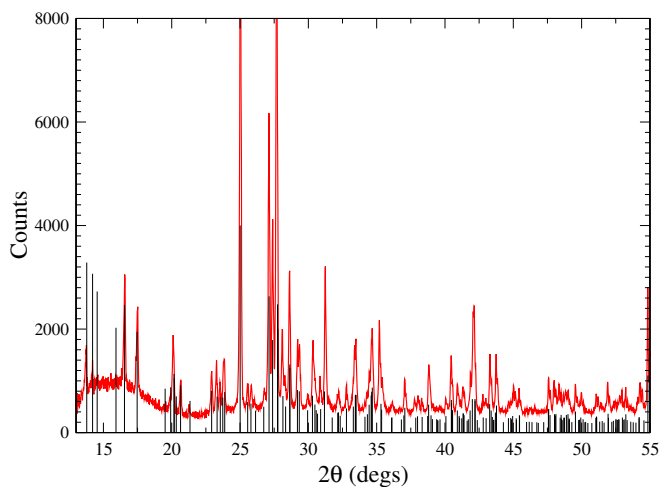


Fig. 1. X-ray diffractogram in our FeVO_4 sample. The bars represent the positions of the expected Bragg peaks in the triclinic system with lattice constants $a=6.709$ Å, $b=8.045$ Å, $c=9.340$ Å, $\alpha=96.43^\circ$, $\beta=106.85^\circ$ and $\gamma=101.49^\circ$ [12].

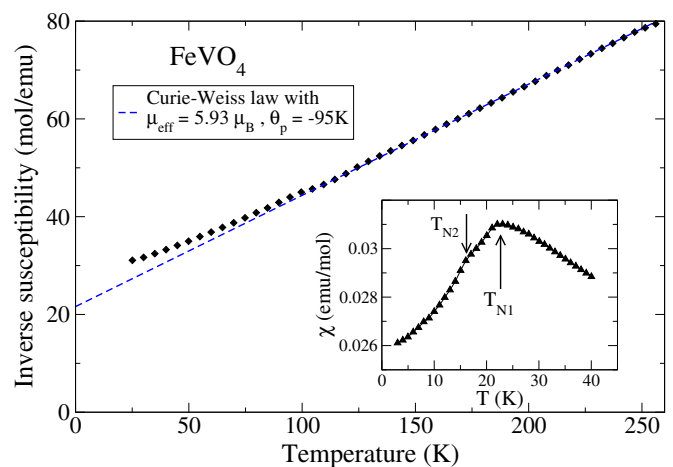


Fig. 2. Thermal variation of the inverse susceptibility in FeVO_4 ; the inset shows the low temperature variation of the susceptibility.

of insulators, the magnetic hyperfine field H_{hf} at the ^{57}Fe nucleus site is proportional to the spontaneous moment with a very good approximation, with a hyperfine constant $C_{\text{hf}} \simeq 11 \text{ T}/\mu_B$. In the following, we shall refer equivalently to the hyperfine field or to the spontaneous moment (except at the lowest temperature, see Section 7). In the magnetically ordered phase, the Mössbauer spectrum associated with a static hyperfine field is a six-line pattern. The quadrupolar hyperfine interaction for $^{57}\text{Fe}^{3+}$ is in general much smaller than the magnetic hyperfine interaction. At first perturbation order, it gives rise to a small lineshift for each line (assuming axial symmetry for the Fe site):

$$\delta e_j = \varepsilon_j \frac{3\Delta E_Q}{4} \left(\cos^2 \theta - \frac{1}{3} \right). \quad (1)$$

In this expression, j is the line index, ε_j is $+1$ for the 2 external lines and -1 for the 4 inner lines, ΔE_Q is the quadrupole parameter whose absolute value can be measured in the paramagnetic phase and θ is the angle between the hyperfine field (or the spontaneous magnetic moment) and the principal axis of the electric field gradient (EFG) tensor at the Fe site. At second perturbation order, the lineshift is

$$\delta^2 e_j = \varepsilon_j \frac{3\Delta E_Q}{4} \left(1 + \beta_j \frac{\Delta E_Q}{h} \sin^2 \theta \right) \cos^2 \theta, \quad (2)$$

where β_j is a coefficient depending on the specific line. The magnetic hyperfine interaction in the excited 14.4 keV nuclear state alone enters here, through the quantity $h = \frac{1}{2} g_n \mu_n H_{\text{hf}}$, where μ_n is the nuclear Bohr magneton and $g_n = -0.10$ is the gyromagnetic ratio of the excited state. In the case of equal moment helical or conical structures, there is no distribution of hyperfine field values and a small spectral effect can arise from the distribution of θ values (if any) associated with the incommensurate structure. It is clear that the first order lineshift yields equivalent broadenings for all the lines. By contrast, the second order shift is different for each line through its dependence on β_j , and its spectral effect consists in inhomogeneous line broadenings. These have been observed in BiFeO_3 [19,20], although in this case it can be shown that they are due to the anisotropy of the hyperfine interaction itself [21] and not to the distribution of θ values associated with the cycloidal spin structure [20]. Inhomogeneous line broadenings due to a helical incommensurate magnetic structure were observed in the langasite compound $\text{Ba}_3\text{NbFe}_3\text{Si}_2\text{O}_{14}$ [22] and in MnGe (doped with Fe) [23].

Much more spectacular effects on the lineshape arise from moment modulated structures, since then the main spectral effect

is due to the distribution of hyperfine field values. This is described in the following, where the small quadrupolar line-shifts have not been considered.

4. Mössbauer lineshapes associated with incommensurate elliptical and sine wave structures

A planar non-collinear elliptical magnetic structure is characterised by the values of the two axes of the ellipse, or by the value of the major axis H_{hf}^{max} and the ratio $y = H_{hf}^{min}/H_{hf}^{max}$. Using the unit vectors **a** and **b** along the principal axes of the ellipse as basis vectors, the hyperfine field writes

$$H_{hf}(\theta) = H_{hf}^{max} (\cos \varphi \mathbf{a} + y \sin \varphi \mathbf{b}), \quad (3)$$

and, in the case of an incommensurate propagation vector, the angle φ is uniformly distributed between 0 and 2π . Then, the hyperfine field distribution at the nucleus site in the interval $H_{hf}^{min} \leq H_{hf} \leq H_{hf}^{max}$ is given by, using $h = H_{hf}/H_{hf}^{max}$:

$$P_{ell}(H_{hf}) \propto \frac{1}{\left| \frac{dH_{hf}(\varphi)}{d\varphi} \right|} = \frac{h}{\sqrt{(h^2 - y^2)(1 - h^2)}}. \quad (4)$$

This distribution is represented on the right panel of Fig. 3, and it can be seen that it diverges at the values H_{hf}^{min} and H_{hf}^{max} ; the corresponding Mössbauer spectrum, shown on the left panel, presents six lines, each of which is split in two peaks, in agreement with the shape of the distribution.

For the case of a collinear sinusoidally modulated structure, the hyperfine field is also given by Eq. (3), but with $y=0$. It is then straightforward to see that the distribution function is given by

$$P_{sin}(H_{hf}) \propto \frac{1}{\sqrt{1 - h^2}}. \quad (5)$$

This function is represented on the right panel of Fig. 4: it is characterised by a divergence for $H_{hf} = H_{hf}^{max}$ and by a nonvanishing weight extending to $H_{hf}=0$. This explains the shape of the corresponding Mössbauer spectrum shown in the left panel of Fig. 4, with a large spectral weight at zero velocity.

5. Mössbauer spectra of FeVO₄

Absorption Mössbauer spectra on the isotope ⁵⁷Fe have been recorded between 1.46 and 25 K, using a commercial Co⁵⁷Rh γ -ray source mounted on a constant acceleration electromagnetic drive. Representative spectra in magnetic phases II (1.46 K) and I (18 K), and in the paramagnetic phase (25 K) are shown in Fig. 5. Our

spectra are in good agreement with those published previously [16,17].

At 25 K, in the paramagnetic phase, the hyperfine quadrupolar interaction alone is present, and the fit must be done with 3 equal weight subspectra corresponding to the 3 crystallographic sites of Fe in FeVO₄. The isomer shifts of these subspectra are around 0.46 (4) mm/s with respect to α -Fe, typical for trivalent Fe in insulators. The quadrupole splittings are close to those given in Ref. [17], namely in decreasing order: 1.08(2), 0.55(2) and 0.24(2)mm/s, corresponding respectively to the green, red and blue subspectrum in Fig. 5 at 25 K. In terms of the components V_{ii} , $i = X, Y, Z$, of the EFG tensor at each site, each quadrupole splitting is worth: $|\Delta E_Q| = |eQV_{ZZ}/2I| \sqrt{1 + \eta^2/3}$, where the asymmetry parameter is: $\eta = (V_{YY} - V_{XX})/V_{ZZ}$. Due to the low site symmetry, η is expected to be non-zero, but it is not possible to obtain independently V_{ZZ} and η . In the following, we will assume $\eta=0$. In addition, the sign of V_{ZZ} cannot be determined from these zero field paramagnetic phase spectra.

In the magnetic phases, the Mössbauer spectra are quite different in phase I and in phase II. They are both six-line patterns, as expected for the magnetic hyperfine interaction of ⁵⁷Fe, but with peculiar shapes: the spectrum in phase II (1.46 K) seems rather complex, with its rightmost peak well resolved, and the spectrum in phase I (18 K) bears a strong resemblance with that shown in Fig. 4.

Therefore, we fitted the spectra in phase II to 3 equal weight subspectra associated with an incommensurate planar elliptical structure, like that shown in Fig. 3. The magnetic structure (for one site) is schematized in Fig. 5 at the right of the 1.46 K spectrum; we have represented the incommensurate elliptical arrangement in projection in the (*ab*) plane around a single Fe atom, the angle φ , which gives the orientation of the hyperfine field in the (*ab*) plane, varying continuously between 0 and 2π . Since the spectra are somewhat asymmetric with respect to zero velocity, quadrupolar effects should be considered for completeness. For this purpose, the knowledge of the electric field gradient (EFG) tensor at the Fe site is in principle required, but the very low symmetry at the Fe sites precludes any *a priori* determination. Since the quadrupolar effects can be considered as a perturbation with respect to the magnetic hyperfine interaction, we used the line shifts given by expression (1) for fitting of the spectra. In a frame where the z-axis is normal to the plane of the ellipse, the principal axis OZ of the EFG tensor is determined by its polar and azimuthal angles θ and Φ . For a given value of φ , the angle θ between the hyperfine field and OZ is such that

$$\cos^2 \theta(\varphi) = \sin^2 \theta \frac{\cos^2(\varphi - \Phi)}{\cos^2 \varphi + y^2 \sin^2 \varphi}, \quad (6)$$

where y is the ratio of the minor to major axis of the ellipse. For a

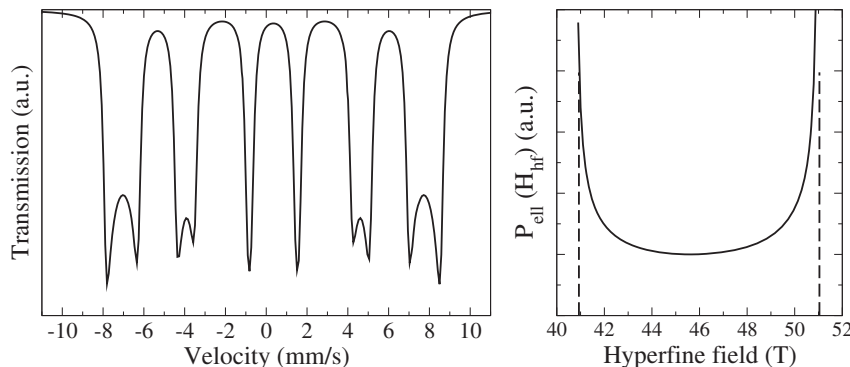


Fig. 3. Right panel: Distribution of hyperfine fields associated with a planar elliptical arrangement of Fe moments with incommensurate propagation vector; the major axis is taken to be 51 T and the ratio of minor to major axes is 0.8; Left panel: Calculated Mössbauer spectrum for such an elliptical structure.

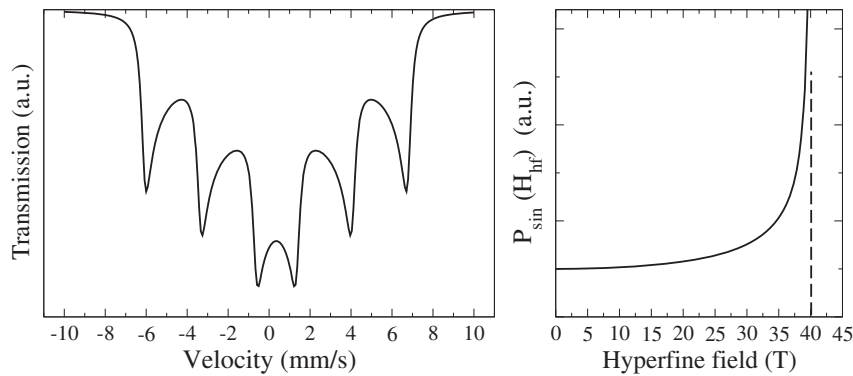


Fig. 4. Right panel: Distribution of hyperfine fields associated with a collinear arrangement of Fe moments with incommensurate sine wave modulation; the maximum field is taken to be 40 T; Left panel: Calculated Mössbauer spectrum for such a sine wave structure.

homogeneous distribution of φ , the angle Φ results mainly in a dephasing of the $\cos^2 \theta$ values and thus has little influence on the spectrum. The fits were performed by letting θ , ΔE_Q , H_{hf}^{max} and y as free parameters for each subspectrum. We find that the obtained quadrupolar parameter values, namely 1.4(3), 0.4(1) and $-0.3(1)$ mm/s have absolute values not far from those measured in the paramagnetic phase, namely 1.08, 0.55 and 0.24 mm/s. However, one must keep in mind the possibility of a lattice distortion occurring at T_N2 , where the ferroelectric order sets in, which could alter the quadrupolar parameter values.

As to the spectra in phase I, they are correctly fitted to 3 equal weight subspectra associated with an incommensurate collinear sine wave structure like that shown in Fig. 4. The magnetic structure (for one site) is schematized in Fig. 5 at the right of the 18 K spectrum, where one period of the sine wave along the propagation vector \mathbf{k} is sketched. The hyperfine field is directed along the principal axis of the ellipse of phase II and thus the angle θ is unique (for a given site). We observe that the effective quadrupole interaction, reflected in the symmetry of the spectra with respect to zero velocity, is very small in phase I. In our

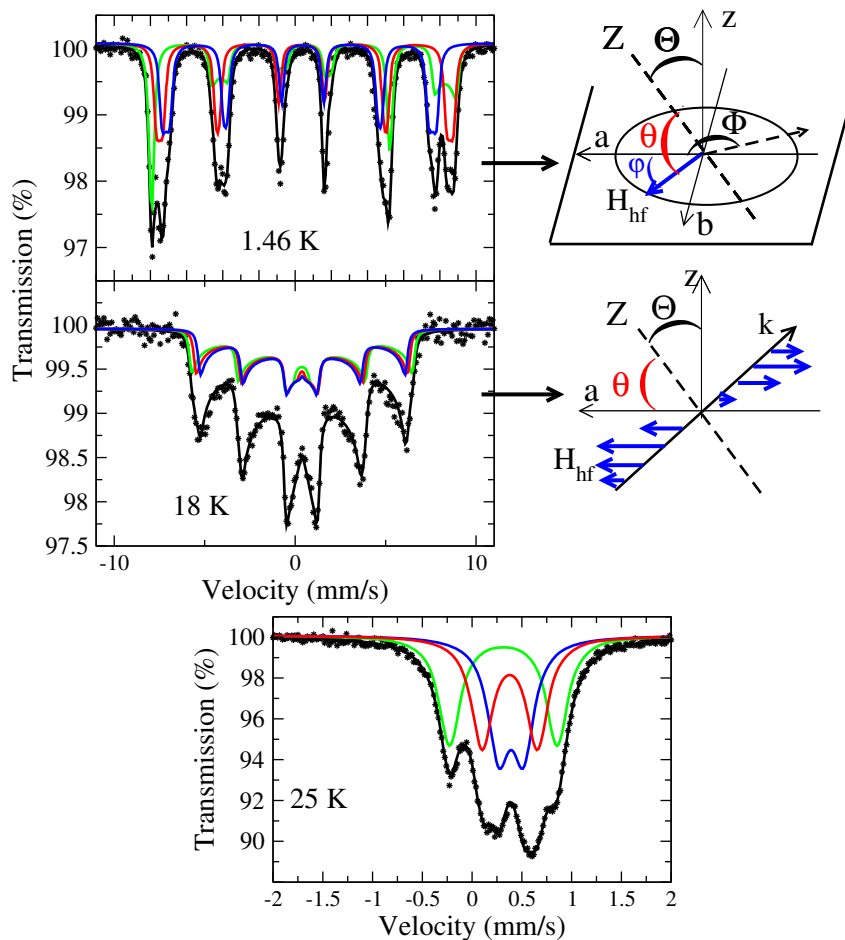


Fig. 5. ^{57}Fe Mössbauer absorption spectra of FeVO_4 at 1.46 K (phase II), 18 K (phase I) and 25 K (paramagnetic phase), the latter with a reduced velocity scale. The lines are fits with 3 subspectra to a planar elliptical magnetic structure (1.46 K), to a collinear sine-wave magnetic structure (18 K) and to a hyperfine quadrupole interaction (25 K). The magnetic structure (for one Fe site) is schematized at the right of each spectrum, with notations as in the text. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this article.)

simulations, we fixed ΔE_Q for the 3 sites at their paramagnetic phase values (1.08, 0.55 and -0.24 mm/s) and, in the presence of the broad distribution of hyperfine fields, we find that the above property is reproduced for relatively large ranges of $\{\theta, \phi\}$ values. It is however not possible to unambiguously determine the actual values for the 3 sites.

6. Thermal variation of the hyperfine fields

Fig. 6 shows the thermal variation of H_{hf}^{max} and y , and Table 1 displays the values obtained at 1.46 K together with a comparison with those derived from neutron diffraction [13] at 2 K. At these base temperatures, our hyperfine field values for the major ellipse axis are in good agreement with those of the major moment axis using a hyperfine constant value ≈ 11.6 T/ μ_B close to the standard value. However, we find that the ratio of minor to major axis is close to 0.93(1) for all 3 sites, larger than the value 0.78(3) obtained by neutron diffraction. As temperature increases and the transition is approached, the ratio y decreases, i.e. the ellipse is getting more and more oblate. This matches well the neutron diffraction finding that the phase II ellipse “merges” into the phase I collinear sine-wave, the major axis of the former becoming the moment direction of the latter. There is good continuity between phase II and phase I, and in phase I at 18 K, the agreement with the neutron values is rather good, with a slightly higher mean hyperfine constant of 12.1 T/ μ_B .

We note that our range of Fe^{3+} moment values at 1.46 K, as well as that in Ref. [13] at 2 K, are at odds with the upper bound of the moment value of $1.95 \mu_B$ derived in the zero field ^{51}V NMR study of Ref. [18] in single crystal $FeVO_4$. This discrepancy could be due to an incorrect estimation of the transferred hyperfine constant at the ^{51}V site in the ordered phase in zero field, which is taken in Ref. [18] to be equal to the high field value in the paramagnetic phase.

7. Discussion

For a spin only ion like Fe^{3+} , no static moment modulation can exist at $T=0$ and therefore, the elliptical structure in phase II should progressively transform into a circular structure as $T \rightarrow 0$.

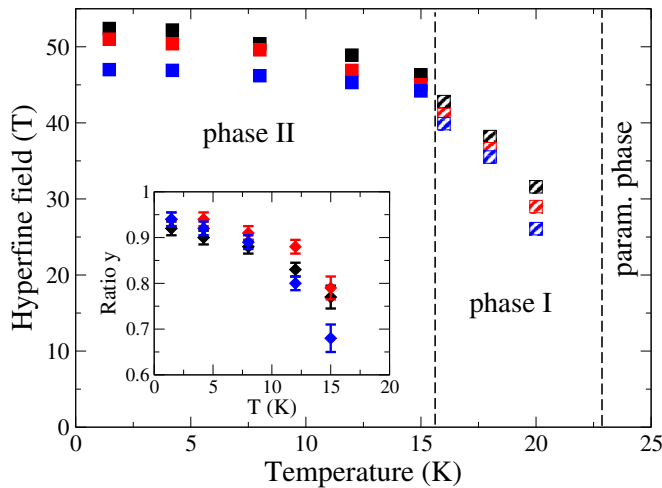


Fig. 6. Thermal variation of the major axis of the hyperfine field ellipse in phase II of $FeVO_4$ (full symbols) and of the maximum hyperfine field in phase I (striped symbols) for the 3 Fe sites. Due to the good statistics of the spectra, the error bar is of the size of the point. The inset shows the thermal variation of the ratio y of the minor to major axes of the hyperfine field ellipse in phase II.

Table 1

For the three sites of Fe in $FeVO_4$: in phase II (planar elliptical structure): maximum hyperfine field H_{hf}^{max} and ratio $y = H_{hf}^{min}/H_{hf}^{max}$ at 1.46 K from the present work, major axis of the elliptical structure m_A and ratio of minor to major moment axes y_n at 2 K according to Ref. [13], deduced hyperfine constant; in phase I (sine wave structure) at 18 K: maximum hyperfine field of the modulation, maximum magnetic moment from Ref. [13], deduced hyperfine constant.

	$H_{hf}^{max}(T)$	y	$m_A (\mu_B)$	y_n	$H_{hf}^{max}/m_A (T/\mu_B)$
Site 1,II	52.4(1)	0.92(1)	4.51(7)	0.81(3)	11.6
Site 2,II	51.0(1)	0.94(1)	4.29(7)	0.79(3)	11.9
Site 3,II	47.0(1)	0.94(1)	4.18(6)	0.76(3)	11.3
Site 1, I	38.2(1)		3.23(5)		11.8
Site 2, I	36.6(1)		3.00(5)		12.1
Site 3, I	35.5(1)		2.86(3)		12.4

In other words, the Fe^{3+} moment on every site should be saturated, i.e. the y ratio should approach 1 as temperature decreases. Our mean y value of 0.93 at 1.46 K is indeed rather close to 1, but it is definitely lower than 1 since we checked that the 1.46 K spectrum cannot be correctly fitted with $y=1$. When extrapolating the $y(T)$ thermal variation to zero temperature (see inset of Fig. 6), one obtains ≈ 0.94 , which is lower than 1. Since this is not allowed for Fe^{3+} , we interpret this deviation from unity as due to the anisotropy of the magnetic hyperfine interaction itself, which should be independent of temperature. This hyperfine anisotropy is documented for $BiFeO_3$ [21], and it can be shown that the angular dependence of the modulus of the hyperfine field can be approximated by

$$H_{hf}(\varphi) = H_{hf}^{\parallel} \cos^2 \varphi + H_{hf}^{\perp} \sin^2 \varphi, \quad (7)$$

where H_{hf}^{\parallel} and H_{hf}^{\perp} are the main components of the elliptical-like trajectory of \mathbf{H}_{hf} . For values of $y_a = H_{hf}^{\parallel}/H_{hf}^{\perp}$ close to 1, it is not possible to distinguish the spectral effects of an elliptical dependence due to a moment modulation (expression (3)) from those due to the anisotropy of the hyperfine interaction (expression (7)). At 1.46 K, the y value reflects solely the hyperfine anisotropy, but at higher temperature, the dominant contribution to the ratio y is the elliptical moment configuration. Regarding the neutron diffraction derived ratio $y_n \sim 0.8$ at 2 K, we have no explanation for such a low value, since it measures directly the ratio of the moment elliptical axes and thus should be much closer to 1.

In another SDW ferroelectric, $FeTe_2O_5Br$, an oblate elliptical incommensurate magnetic structure of Fe^{3+} moments has been observed to persist down to 0.053 K [24,25], with a quite small y ratio of 0.37, seemingly violating the “single-valued moment” rule for Fe^{3+} as $T \rightarrow 0$. However, a fluctuating disordered moment component has been inferred from μSR measurements down to very low temperature, which should restore a single “static” Fe^{3+} moment on each site as $T \rightarrow 0$. This is confirmed by our Mössbauer data in $FeTe_2O_5Br$ [26] at 4.2 K, which shows a well resolved magnetic hyperfine spectrum with a single hyperfine field of ≈ 44 T. This is in good agreement with the moment value of $4 \mu_B$ quoted in Ref. [24] using the standard hyperfine constant of 11 T/ μ_B . The hyperfine Larmor period associated with the ^{57}Fe magnetic hyperfine interaction is $\tau_M \sim 10^{-8}$ s, so the fluctuations of the disordered spin component in $FeTe_2O_5Br$ must be rather slow, with a characteristic time longer than τ_M . Therefore, the persistent spin dynamics at play in $FeTe_2O_5Br$ down to the lowest temperature does not violate the “single-valued moment” rule for Fe^{3+} as $T \rightarrow 0$.

8. Conclusion

Magnetic susceptibility data and ^{57}Fe absorption Mössbauer spectra have been recorded for FeVO_4 in the two incommensurate magnetic phases, with transition temperatures 15.7 and 23 K. The spectra are in very good agreement with the neutron diffraction results for this compound and represent a good illustration of the spectral shapes associated with magnetic phases with incommensurate moment modulations, in the present case planar elliptic and collinear sine-wave. We observe that the ratio of minor to major elliptical hyperfine field axes is not exactly unity as $T \rightarrow 0$, which we interpret as due to the anisotropy of the magnetic hyperfine interaction since no moment modulation can exist for Fe^{3+} as $T \rightarrow 0$. Mössbauer spectroscopy is an important complementary method of neutron diffraction for studying Fe (or Sn) containing moment modulated magnetic phases, yielding rather precise values for the characteristics of the magnetic structures and often leading to a more thorough understanding of the system under study.

References

- [1] G.A. Smolenskii, I.E. Chupis, *Sov. Phys. Usp.* 25 (1982) 475.
- [2] M. Mostovoy, *Phys. Rev. Lett.* 96 (2006) 067601.
- [3] Y. Tokura, S. Seki, N. Nagaosa, *Rep. Prog. Phys.* 77 (2014) 076501.
- [4] R.T. Smith, G.D. Achenbach, R. Gerson, W.J. James, *J. Appl. Phys.* 39 (1968) 70.
- [5] I. Sosnowska, T. Peterlin-Neumaier, E. Steichele, *J. Phys. C* 15 (1982) 4835.
- [6] A. Kadomtseva, A. Zvezdin, Y. Popov, A. Pyatakov, G. Vorob'ev, *JETP Lett.* 79 (2004) 571.
- [7] T. Kimura, S. Ishihara, H. Shintani, T. Arima, K.T. Takahashi, K. Ishizaka, Y. Tokura, *Phys. Rev. B* 68 (2003) 060403 (R).
- [8] T. Kimura, T. Goto, H. Shintani, K. Ishizaka, T. Arima, Y. Tokura, *Nature (London)* 426 (2003) 55.
- [9] N. Hur, S. Park, P.A. Sharma, J.S. Ahn, S. Guha, S.W. Cheong, *Nature (London)* 429 (2004) 392.
- [10] A. Yoshimori, *J. Phys. Soc. Jpn.* 14 (1959) 807.
- [11] I. Dzyaloshinski, *J. Exp. Theor. Phys.* 19 (1964) 960 [*Zh. Eksp. Teor. Fiz.* 46 (1964) 1420].
- [12] Z. He, J.-I. Yamaura, Y. Ueda, *J. Solid State Chem.* 181 (2008) 2346.
- [13] A. Daoud-Aladine, B. Kundys, C. Martin, P.G. Radaelli, P.J. Brown, C. Simon, L. C. Chapon, *Phys. Rev. B* 80 (2009) 220402 (R).
- [14] B. Kundys, C. Martin, C. Simon, *Phys. Rev. B* 80 (2009) 172103.
- [15] M. Kenzelmann, A.B. Harris, S. Jonas, C. Broholm, J. Schefer, S.B. Kim, C. L. Zhang, S.W. Cheong, O.P. Vajk, J.W. Lynn, *Phys. Rev. Lett.* 95 (2005) 087206.
- [16] L.M. Levinson, B.M. Wanklyn, *J. Solid State Chem.* 3 (1971) 131.
- [17] B. Robertson, E. Kostiner, *J. Solid State Chem.* 4 (1972) 29.
- [18] J. Zhang, L. Ma, J. Dai, Y.P. Zhang, Zhangzhen He, B. Normand, Weiqiang Yu, *Phys. Rev. B* 89 (2014) 174412.
- [19] A. Palewicz, T. Szumiata, R. Przenioslo, I. Sosnowska, I. Margiolaki, *Solid State Commun.* 140 (2006) 359.
- [20] D. Lebeugle, D. Colson, A. Forget, M. Viret, P. Bonville, J.F. Marucco, S. Fusil, *Phys. Rev. B* 76 (2007) 024116.
- [21] A.V. Zalesky, A. Frolov, T.A. Khimich, A.A. Bush, V.S. Pokatilov, A.K. Zvezdin, *Europhys. Lett.* 50 (2000) 547.
- [22] K. Marty, P. Bordet, V. Simonet, M. Loire, R. Ballou, C. Darie, J. Kljun, P. Bonville, O. Isnard, P. Lejay, B. Zawilski, C. Simon, *Phys. Rev. B* 81 (2010) 054416.
- [23] M. Deutsch, P. Bonville, A.V. Tsvyashchenko, L.N. Fomicheva, F. Porcher, F. Damay, S. Petit, I. Mirebeau, *Phys. Rev. B* 90 (2014) 144401.
- [24] M. Pregelj, O. Zaharko, A. Zorko, Z. Kutnjak, P. Jeglić, P.J. Brown, M. Jagličić, Z. Jagličić, H. Berger, D. Aržon, *Phys. Rev. Lett.* 103 (2009) 147202.
- [25] M. Pregelj, A. Zorko, O. Zaharko, D. Aržon, M. Komelj, A.D. Hillier, H. Berger, *Phys. Rev. Lett.* 109 (2012) 227202.
- [26] D. Colson, A. Forget, P. Bonville, unpublished results.